

## Embedding

To fully appreciate a sphere we embed it in 3D flat space

At least 5D is required to represent 4D spacetime

Alternative - take 2D slice of 4D spacetime  
+ embed slice in 3D

Recall 2D space  $ds^2 = a^2(d\theta^2 + \sin^2\theta d\phi^2)$

and 3D flat space  $ds_{3\text{flat}}^2 = d\rho^2 + \rho^2(d\chi^2 + \sin^2\chi d\varphi^2)$

in general, write  $\rho = \rho(\theta, \phi)$ ;  $\chi = \chi(\theta, \phi)$ ;  $\varphi = \varphi(\theta, \phi)$

an example  $ds^2 = -dt^2 + dr^2 + (b^2 + r^2)(d\theta^2 + \sin^2\theta d\phi^2)$

$r$  can be pos or neg

- metric is curved
- at fixed  $r, t$  geometry is that of a 2-sphere with radius  $r^2 + b^2$
- for  $r \rightarrow \pm\infty$   $ds^2 \rightarrow ds_{\text{flat}}^2$  asymptotically flat

- spacetime is time-independent - focus on space part

$$dS_{3D}^2 = dr^2 + (b^2 + r^2)(d\theta^2 + \sin^2\theta d\phi^2)$$

- spherical symmetry means we can take any slice  
e.g.  $\theta = \pi/2$  - equatorial plane

$$d\Sigma^2 = dr^2 + (b^2 + r^2)d\phi^2$$

embed in 3D flat space

$$dS_{3Dflat}^2 = d\rho^2 + \rho^2 d\varphi^2 + dz^2$$

use  $\varphi = \phi$  ;  $\rho = \rho(r)$  ;  $z = z(r)$

$$dS_{3Dflat}^2 \text{ on surface} = \left( \left( \frac{d\rho}{dr} \right)^2 + \left( \frac{dz}{dr} \right)^2 \right) dr^2 + \rho^2(r) d\phi^2$$

so  $\left( \frac{d\rho}{dr} \right)^2 + \left( \frac{dz}{dr} \right)^2 = 1$        $\rho^2 = r^2 + b^2$

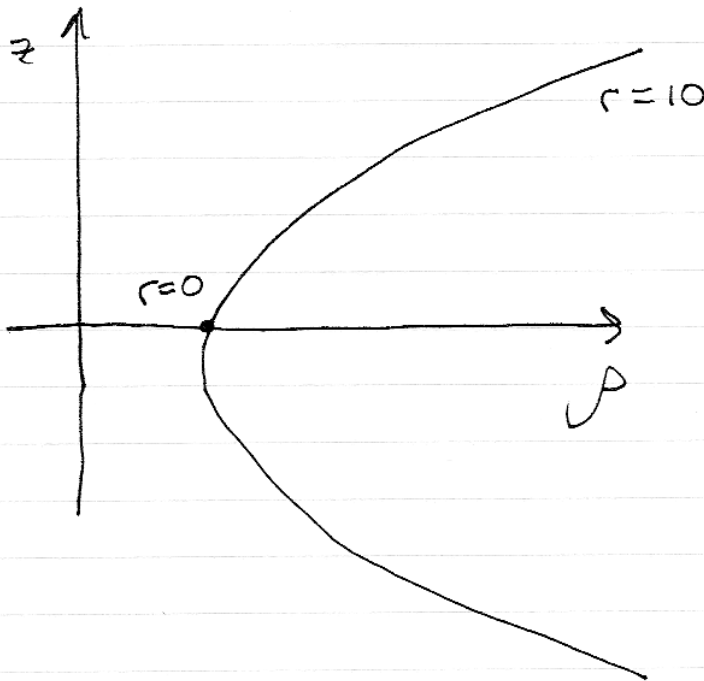
$$\Rightarrow \rho d\rho = r dr \quad \Leftrightarrow \quad \frac{dz}{dr} = \left(1 - \frac{r^2}{\rho^2}\right)^{1/2}$$

$$= \frac{b}{(r^2 + b^2)^{1/2}}$$

$$z = b \int \frac{dr}{(r^2 + b^2)^{1/2}} = b \sinh^{-1}(r/b)$$

$$\text{or } r = b \sinh(z/b)$$

$$\rho = (r^2 + b^2)^{1/2} = b \cosh(z/b)$$



## Vectors

1st year definition - quantity with magnitude + direction  
obeys rules of addition, scalar products ...

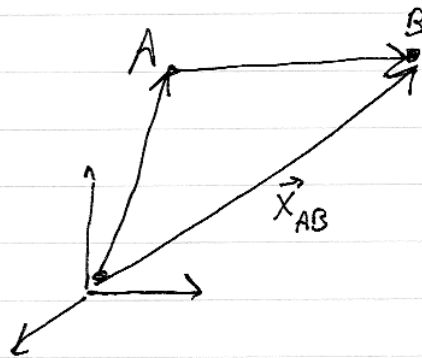
e.g. displacement, velocity, acceleration

of these, displacement is different - non-local

$$\vec{x}_{AB} = \vec{x}_B - \vec{x}_A$$

$$\vec{v} = \left. \frac{d\vec{x}}{dt} \right|_{x=x_p, t}$$

$$\vec{a} = \left. \frac{d\vec{v}}{dt} \right|_{x=x_p, t}$$



In general relativity, only local vectors make sense

Can write vectors in terms of basis

$$\underline{a}(x) = a^\alpha(x) \underline{e}_\alpha(x) \quad \alpha \text{ labels basis vectors}$$

$$\left( \underline{e}_\alpha \right)^\beta \text{ is } \beta \text{ comp of } \underline{e}_\alpha$$

Put another way

$$u^\alpha = \frac{dx^\alpha}{d\tau} \quad \text{only means something if we specify basis}$$

$$\text{We know } \underline{u} \cdot \underline{u} = -1 = g_{\alpha\beta} \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau}$$

$$= u^\alpha \underline{e}_\alpha \cdot u^\beta \underline{e}_\beta$$

$$\Rightarrow \underline{e}_\alpha \cdot \underline{e}_\beta = g_{\alpha\beta} \quad \text{by definition}$$

Scalar product  $\underline{a} \cdot \underline{b} = \underline{e}_\alpha \cdot \underline{e}_\beta a^\alpha b^\beta$

2 choices for basis

• orthonormal  $\underline{e}_\alpha \cdot \underline{e}_\beta = \eta_{\alpha\beta}$

i.e., 4 unit vectors

think observer with  $\underline{u}_{\text{obs}} = \underline{e}_t = (1, 0, 0, 0)$  rest frame

+ 3 orthogonal unit vectors

$\underline{p}$  4-momentum of part. in lab

$$p^{\hat{t}} = - \underline{p} \cdot \underline{u}_{\text{obs}} = E = \text{energy}$$

• coordinate basis  $\underline{e}_\alpha \cdot \underline{e}_\beta = g_{\alpha\beta}$

$$\underline{a} \cdot \underline{b} = g_{\alpha\beta} a^\alpha b^\beta$$

e.g.  $\underline{u} \cdot \underline{u} = g_{\alpha\beta} u^\alpha u^\beta \equiv g_{\alpha\beta} \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau}$   $u^\alpha \equiv \frac{dx^\alpha}{d\tau}$  in coordinate basis

$$= \frac{g_{\alpha\beta} dx^\alpha dx^\beta}{(d\tau)^2} = -1$$